## **Engineering Notes**

## Three-Fourths $T_{\text{max}}$ Rule for Minimum-Area Space Radiators

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ETAILED thermodynamic calculations of space power systems can be very laborious. A number of short cuts have been devised, but such computational aids must be used with caution, especially when comparing systems involving different cycle characteristics and/or different working fluids. For example, Moffitt and Klag¹ have demonstrated that if one assumes a Carnot cycle or a saturated Rankine cycle (Fig. 1a) and a radiator that is an abstraction of real radiators (i.e., a radiator without tubes), then the minimum area radiator temperature will be three-fourths of the maximum cycle temperature (in degrees absolute). This  $\frac{3}{4}$   $T_{\text{max}}$  relationship is often used as a rule of thumb. However, it is apparent from Fig. 1a that the saturated Rankine cycle shown is not grossly different from a Carnot cycle, so that the  $\frac{3}{4}$   $T_{\text{max}}$  rule should be applicable. This cycle is quite popular, and working fluids suggested for it include mercury,2 potassium,3 rubidium,4 and, in an unusual version (which expands from saturated to superheated), biphenyl.<sup>5</sup> Even with such cycles, however, efficiency effects tend to change the value for minimum-area radiator temperature, as has been pointed out by Szego.

Finally, there are systems in which the cycle departs so markedly from the Carnot form that it is not even approximately correct to assume that the  $\frac{3}{4}$   $T_{\rm max}$  relationship gives minimum radiator areas. For example, water, which has been recently proposed, requires a significant amount of superheat in order to achieve competitive weight powerplants. The temperature-entropy diagrams of such cycles can differ markedly in morphology from that of the Carnot cycle, as suggested by Fig. 1b. For such systems, the  $\frac{3}{4}$   $T_{\rm max}$  relationship is definitely not valid.

As can easily be seen by examination, the assumption that the useful work is proportional to the differences between  $T_{\rm max}$  and  $T_{\rm rad}$ , as must be assumed to derive the three-fourths rule, must be quite erroneous. To demonstrate this, assume isentropic expansion, as well as a fluid with a very steep slope for the constant pressure line in the super-

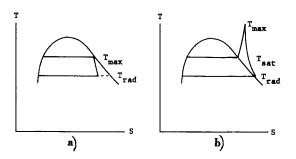


Fig. 1 Temperature entropy diagrams of Rankine cycles: a) expanding from saturated vapor, b) expanding from superheated vapor.

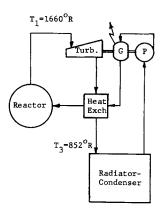


Fig. 2 60 ekw single-loop steam system.<sup>7</sup> Over-all efficiency, 14%; specific weight, 52 lb/ekw; radiator specific weight, 18 lb/ ekw; no auxiliary cooling loops.

heat region (i.e., very low specific heat). For such a cycle, the area under the superheated "spike" will approach zero, and the temperature that could then be used in the three-fourths assumption with some degree of accuracy would be that of the saturated vapor  $T_{\rm sat}$  rather than the top temperature  $T_{\rm max}$ .

It is interesting to compare results of these two assumptions with a case (Fig. 2) which has been analyzed in detail. (This 60-ekw powerplant for a synchronous orbit satellite uses steam, with a top temperature of 1660°R, a saturation temperature of 1027°R, and a regenerative heat exchanger between turbine exhaust and reactor inlet.) The results are tabulated in Table 1. Note that the correct temperature for a minimum area radiator is 852°R, which is only 51% of the top temperature.

Table 1 Comparison of values and ratios for three approaches for obtaining the minimum area radiator temperature

Approaches	Data		
Relationship used	$\frac{3}{4} T_{\text{max}}$	$\frac{3}{4} T_{\mathrm{sat}}$	Actual min area, T <sub>rad</sub>
$T_{ m rad}$ , from relationship °R Ratio $T_{ m rad}/T_{ m max}$	$1245 \\ 0.75$	$770 \\ 0.46$	$852 \\ 0.51$

## References

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